

Lecture 1 Early stage of quantum physics

Monday, April 29, 2024

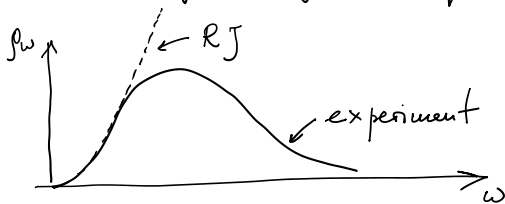
(1.1) Thermal e/m radiation and ultraviolet catastrophe

End of XIX century

- All bodies with $T \neq 0$ are sources of e/m radiation
- Thermal radiation: the body is in thermal equilibrium with all surrounding bodies $\Rightarrow$ determined by T
- Black body (black body cavity): ideal absorber/emitter of thermal radiation (no radiation passing through or being reflected)
- Experiments on black body radiation

$\rightarrow$ energy density $u = \frac{U}{V} = \int_0^\infty \rho_\omega d\omega$

$\uparrow$ spectral density of e/m radiation



$\rightarrow$ exper: $u = \sigma T^4$ [Stefan-Boltzmann]

$\rightarrow$ exper: $\omega_{max} \propto T$ [Wien's displacement]

- Classical theory of e/m radiation

$$\rho_\omega = \frac{\omega^2}{\pi^2 c^3} k_B T \quad [\text{Rayleigh-Jeans}]$$

$u \sim \infty$??? Ultraviolet catastrophe

(1.2) Planck hypothesis and Planck constant

Nov 1900 Planck RJ $\Rightarrow \rho_\omega = \frac{\omega^2}{\pi^2 c^3} \langle E \rangle$

$\uparrow$ mean energy of oscillator/standing e/m wave in cavity
number of oscillators in RJ theory

Planck's revolutionary hypothesis: energy of oscillator changes with discrete quanta of energy $h\nu$. $h = \frac{h}{2\pi}$ Planck constant

$$\langle E \rangle = \frac{1}{2} h\nu + \frac{h\nu}{e^{h\nu/k_B T} - 1} = \frac{1}{2} h\nu \coth \frac{h\nu}{2k_B T}$$

quantum Gibbs

Modern result

$$E_n = h\nu \left(n + \frac{1}{2}\right)$$

$$\rho_\omega = \left(\rho_\omega^{(0)}\right) + \rho_\omega^T$$

forget about $\rho_\omega^{(0)}$

$$\boxed{\rho_\omega^T = \frac{\omega^2}{\pi^2 c^3} \frac{h\nu}{e^{h\nu/k_B T} - 1}}$$

[Planck formula]

Remark 1. Angular frequency $\omega = 2\pi\nu$

$$\rho_\omega d\omega = \rho_\nu d\nu \Rightarrow \rho_\nu = \rho_\omega \Rightarrow \frac{d\omega}{d\nu} = \frac{8\pi h \nu^3}{c^3} \frac{1}{e^{h\nu/k_B T} - 1}$$

Remark 2. At $\frac{h\nu}{k_B T} \ll 1 \Rightarrow \text{Planck} \rightarrow \text{RJ}$

Remark 3. $u = \int_0^\infty \rho_\omega^T d\omega = \alpha T^4 [\text{S-B}]$ exp. solve $\alpha \Rightarrow \text{velocity}$

(1.3) Concept of photon

(a) photoelectric effect: emission of electrons from metal surface caused by light — major puzzle in physics in early 1900s

Experimental laws of phot. effect:

- 1) number of electrons emitted per unit time $\propto$ light intensity:
- 2) kinetic energy of emitted electrons does not on light intensity:
- 3) effect possible if $\omega > \omega_{cr}$
kinetic energy of electrons $\propto \omega - \omega_{cr}$

E/m theory unable to explain 2) and 3)

Einstein 1905 with energy balance equation

$$h\nu = A + \frac{mv^2}{2}$$

$\uparrow$ energy of light quanta (photon) $\uparrow$ work function: energy required to remove electron from the surface $\leftarrow$ kinetic energy of electron

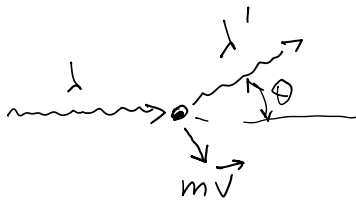
$$A = h\nu_{cr} \Rightarrow \frac{mv^2}{2} = h(\omega - \omega_{cr}) \quad \# 3)$$

(b) Compton effect : interaction of X-rays (Röntgen rays, high-energy e/m radiation) with matter

Classical e/m theory (Thomson scattering)

Compton (1922-23) : spectrum of scattered radiation λ, ω (incident radiation) AND $\lambda', \omega', \omega' < \omega, \lambda' > \lambda$

contradicts classical e/m theory



$$\Delta\lambda = \lambda' - \lambda = 2\lambda_e \sin^2(\theta/2)$$

$$\lambda_e \sim 2.4 \cdot 10^{-10} \text{ cm Compton constant}$$

Explanation : Momentum of photon $\vec{p} = \hbar \vec{k}$

$$k = \frac{2\pi}{\lambda} = \frac{\omega}{c}$$

Two conservation laws

$$\hbar\omega = \hbar\omega' + \frac{m_e v^2}{2}$$

$$\hbar\vec{k} = \hbar\vec{k}' + m_e \vec{v}$$

$$\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

Conclusion : E/m radiation behave like a wave (interference and diffraction experiments), but also have a particle nature (photoelectric and Compton effects)

(1.4) Nuclear model of atom

(a) Two basic models

known at the beginning of XX century :

- existence of positive ions

- electron

- atom is complex electrical system $\sim 10^{-8} \text{ cm}$;

- model ?

- Thomson's model (plum pudding model)
Atoms are uniform spheres of positively charged matters in which electrons are embedded;
- Planetary model:
 - sun $\leftrightarrow$ positive nucleus with Ze , Z serial No of an element in periodic table
 - planets $\leftrightarrow$ electrons.

(b) Rutherford experiments on scattering α -particles (helium nuclei = 2 protons, 2 neutrons) $\Rightarrow$

- $\Rightarrow$ planetary model
- upper bound of nucleus size
 $r \sim 10^{-13}$ cm $\ll$ atom size $\sim 10^{-8}$ cm

- problem of planetary model:
 - angular velocity of electron
 $m\omega^2 r = \frac{e^2}{r^2} \Rightarrow \omega = \sqrt{\frac{e^2}{mr^3}}$

- classical e/m theory: energy of electron $\rightarrow$
 $\rightarrow$ energy e/m radiation

$$\frac{dE}{dt} = -gE \quad g = \frac{2e^2\omega^2}{3mc^3} \quad \text{radiation friction}$$

$\sim 10^8 \text{ sec}^{-1}$

Conclusion: after 10^{-8} sec electron "falls down" on the nucleus $\Rightarrow$ atom is unstable system ???

(1.5) Bohr theory

(a) Bohr's postulates

- Atom exists for a long time in certain stationary states with discrete spectrum of energies

$$E_1, E_2, \dots$$

- Transition from one stationary state with E_n to another $E_m < E_n$ is accompanied by emission of light quanta

$$h\nu_{nm} = E_n - E_m$$

(b) Quantization rules for circular orbits of electrons in atom

- Angular momentum $m_e v_e r = n h$ (*)
 n : principal quantum number, $n=1, 2, 3$

- Radius of orbit from Newton law

$$\frac{m v^2}{r} = \frac{Z e^2}{r^2} \Rightarrow v^2 = \frac{Z e^2}{m_e r} = \frac{n^2 h^2}{m_e^2 r^2} \Rightarrow$$

$$\Rightarrow \boxed{r = \frac{h^2 n^2}{Z m_e e^2}} \quad (**)$$

- 1st Bohr orbit (of hydrogen atom), $n=1$, $Z=1$

$$r_0 \equiv a_0 = \frac{h^2}{m_e e^2} \approx 0.5 \cdot 10^{-8} \text{ cm}$$

- energy of electron in atom at n -th level

$$E_n = E_{kin} + E_{pot} = - \frac{Z e^2}{2r} = - \frac{Z e^2}{2} \frac{Z m_e e^2}{h^2 n^2} \Rightarrow$$

$$\frac{m_e v^2}{2} - \frac{Z e^2}{r} \quad (**)$$

$$(*) \rightarrow \frac{Z e^2}{2r}$$

$$\Rightarrow \boxed{E_n = - \frac{m_e Z^2 e^4}{2 h^2 n^2}} \quad (***)$$

$n=1$
 energy of ground state

- Remarks
- 1) $\min E_n$ at $n=1$
 - 2) $n \rightarrow \infty$ $E_n \rightarrow 0$

(c) Two more quantum numbers (Sommerfeld, Debye)

- q.n. l : determines the shape of electron's orbital
 $l = 0, 1, \dots, n$

- magnetic q.n. m : orientation in space of particular orbital

$m\hbar$ = component of atom's angular momentum along the magnetic field

$$-l \leq m \leq l, \quad 2l+1 \text{ values of } m \Rightarrow$$

$\Rightarrow$ number of allowed directions of angular momentum vector in magnetic field:
space quantization of angular momentum

(d) Summary of early amendments to Bohr theory

- Energy of electron in atom is determined by 3 q.n.
- Transition from one energy state to another involves change of one or more q.n.
- In transition electron absorbs or emits light of particular frequency

$$E_{\text{final}} - E_{\text{initial}} = h\nu = h\omega$$

- However: atom is placed in magnetic field $\Rightarrow$
 $\Rightarrow$ multiplicity of spectra (all the observed frequencies) cannot be explained by space quantization rules alone (with n, l, m)

Sommerfeld (1920): Another quantum number (inner q.n.)?

(1.6) Birth of "spin".

- Kronig (1925): in addition to orbital angular momentum $\Rightarrow$ angular momentum caused by spinning of electron about its own axis

Kronig postulated ang mom $= \frac{h}{2}$ (why?)

- Kronig able to explain the multiplicity of spectra, but: rotation rate is 130 times of $c = 3 \cdot 10^8$ m/sec $\Rightarrow$

$\Rightarrow$ Kronig never published his idea

- Uhlenbeck & Goudsmit: the same spinning idea in *Naturwissenschaften*

- Kronig: letter in *Nature* criticizing U-G

- U-G: publish 2nd paper in *Nature*
1st paper has a discrepancy of the factor of 2 with experimental result

- Thomas (1926): Factor of 2 correction

BUT

Q1. Why the spin angular momentum is quantized?

Q2. Why the idea cannot be reconciled with special theory of relativity?

Chapter 2. Schrödinger Equation

(2.1) DeBroglie waves, Wave properties of particles

- E/m radiation: first wave-like, then particle-like properties (corpuscular properties of light)

Photon $E = h\omega$, $p = h\vec{k}$ (1)

- Matter: first particle-like then wave-like properties

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deBroglie hypothesis (1923): '(1)' are valid for
particles with non-zero rest mass

Wave-like characteristics

$$\omega = \frac{E}{\hbar}, \quad \lambda_B = \frac{2\pi}{k} = \frac{2\pi\hbar}{p} = \frac{2\pi\hbar}{mv} = \frac{h}{mv} \quad (2)$$

↑
DeBroglie wavelength

⇒ wave field for particles characterized by
wave function $\psi(\vec{r}, t)$

- probabilistic interpretation

$$|\psi(\vec{r}, t)|^2 \frac{d\vec{r}}{V} = \text{Pr} \{ \vec{r} \leq \vec{r}(t) < \vec{r} + d\vec{r} \} \quad (3)$$

$$\int |\psi(\vec{r}, t)|^2 \frac{d\vec{r}}{V} = 1$$

- Monochromatic deBroglie wave (E and $\vec{p}$ fixed)
Analogue of $e^{-i\omega t + i\vec{k}\vec{r}}$

$$\psi(\vec{r}, t) = C \exp\left(-\frac{iEt}{\hbar} + \frac{i\vec{p}\vec{r}}{\hbar}\right) \quad (4)$$

$$|\psi|^2 = \psi\psi^* = C^2$$

Remark. Normalization If (3) ⇒ $C = 1$

$$\text{If } \int |\psi(\vec{r}, t)|^2 d\vec{r} = 1 \Rightarrow C = \frac{1}{\sqrt{V}} \quad (5)$$

Remind Bohr quantization rule for angular momentum

$$\underbrace{m_e v r_n}_{\uparrow \frac{h}{\lambda_B}} = n\hbar \Rightarrow \boxed{2\pi r_n = n \lambda_B} \quad (6)$$

stationary orbits: integer number of deBroglie waves

Experiments fully confirm de Broglie hypothesis about wave properties of the particles with non-zero mass.

(2.2) Wave equation for particles - Schrödinger eq

(a) Equation for free particle

$$E = \frac{p^2}{2m} \Rightarrow \boxed{\hbar\omega = \frac{\hbar^2 k^2}{2m}} \quad (7)$$

example of quantum dispersion relation

$$\Downarrow$$

$$\boxed{i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t} = -\frac{\hbar^2}{2m} \Delta \psi(\vec{r}, t)} \quad (8)$$

Check $\psi(\vec{r}, t) \sim \exp(-i\omega t + i\vec{k}\vec{r})$

(b) Momentum and Hamiltonian operators

$$\hat{H} = \hat{K} = \frac{\hat{p}^2}{2m} = -\frac{\hbar^2}{2m} \Delta_r$$

Momentum operator $\hat{p} = -\frac{i}{\hbar} \frac{\partial}{\partial \vec{r}} = -\frac{i}{\hbar} \nabla_r$, from (8)

$$\boxed{i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t} = \hat{H} \psi(\vec{r}, t)} \quad (9)$$

• Next step: Eq. (9) is valid in general case

Example 1. Particle in potential field $U(\vec{r}, t)$

$$\hat{H} = -\frac{\hbar^2}{2m} \Delta_r + U(\vec{r}, t) = \frac{\hat{p}^2}{2m} + U(\vec{r}, t)$$

Example 2. N-particle system

$$\hat{H} = \sum_{1 \leq i \leq N} \left(-\frac{\hbar^2}{2m} \Delta_{\vec{r}_i} + U(\vec{r}_i, t) \right) + \sum_{1 \leq i < j \leq N} \Phi(|\vec{r}_i - \vec{r}_j|)$$

(c) Stationary solution

Go back to (9), seek solution for the time independent

$$\Psi(\vec{r}, t) = \exp\left(-\frac{iEt}{\hbar}\right) \psi_E(\vec{r})$$

$$\hat{H} \psi_E(\vec{r}) = E \psi_E(\vec{r})$$

$\Downarrow$
 $\psi_E(\vec{r})$ eigenfunction of Hamilton, E eigenvalue

Dual role $\hat{H}$

1. $i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$

time evolution of the system considered

2. $\hat{H} \psi_E = E \psi_E$

energies and eigenfunctions
of stationary states

Q. How to include spin in Schrödinger eq?